

Fundamentals of Magnetic Island Theory in Tokamaks

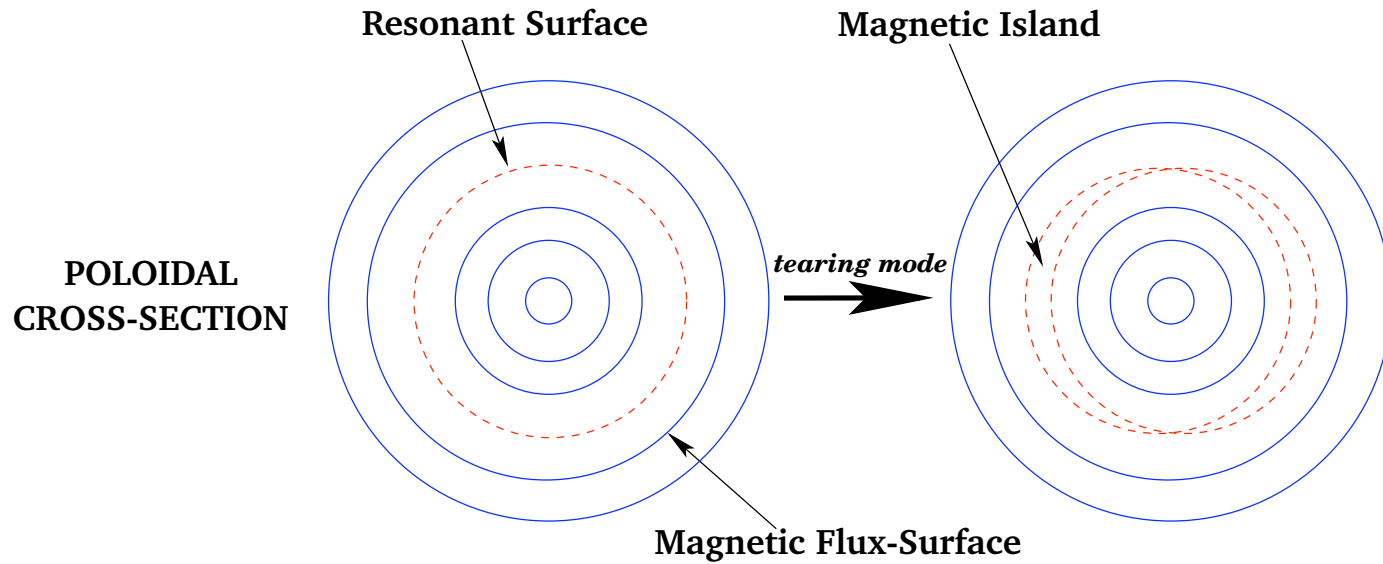
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<http://farside.ph.utexas.edu/talks/talks.html>

Magnetic Islands



- $\vec{k} \cdot \vec{B} = 0$

rational magnetic surfaces
- ...

Need for Magnetic Island Theory

- gnc s n fr t n ssc t th *nonlinear* ph s f
 t r ng gr th e h nr s n th c s
 gr t r th n n r yr th t r t n s rfc
- n ry h t p s s f n n r n y t s n r y r s
 s th n th t t r ng r y n n n n r r g h n r s t
 t^c t
- L n r t r ng th ry r g y r r n R q r n n n r
 gnc s n th ryt p n p r n s r t ns

MHD Theory

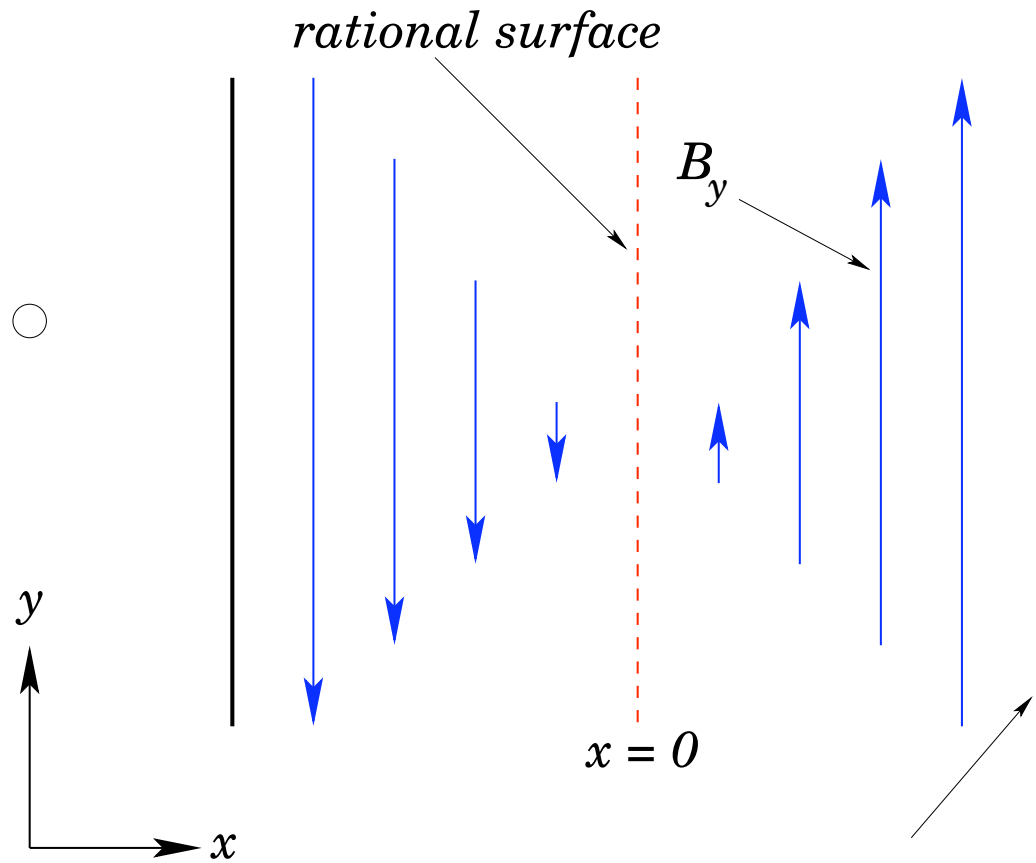
- ring structure consists of the central column of plasma with a central axis of symmetry. The theory is based on the single-fluid theory.

- The theory is based on the magnetohydrodynamical approximation^a for the central column of plasma. The single-fluid theory is used.

- The slab approximation is used to describe the plasma.

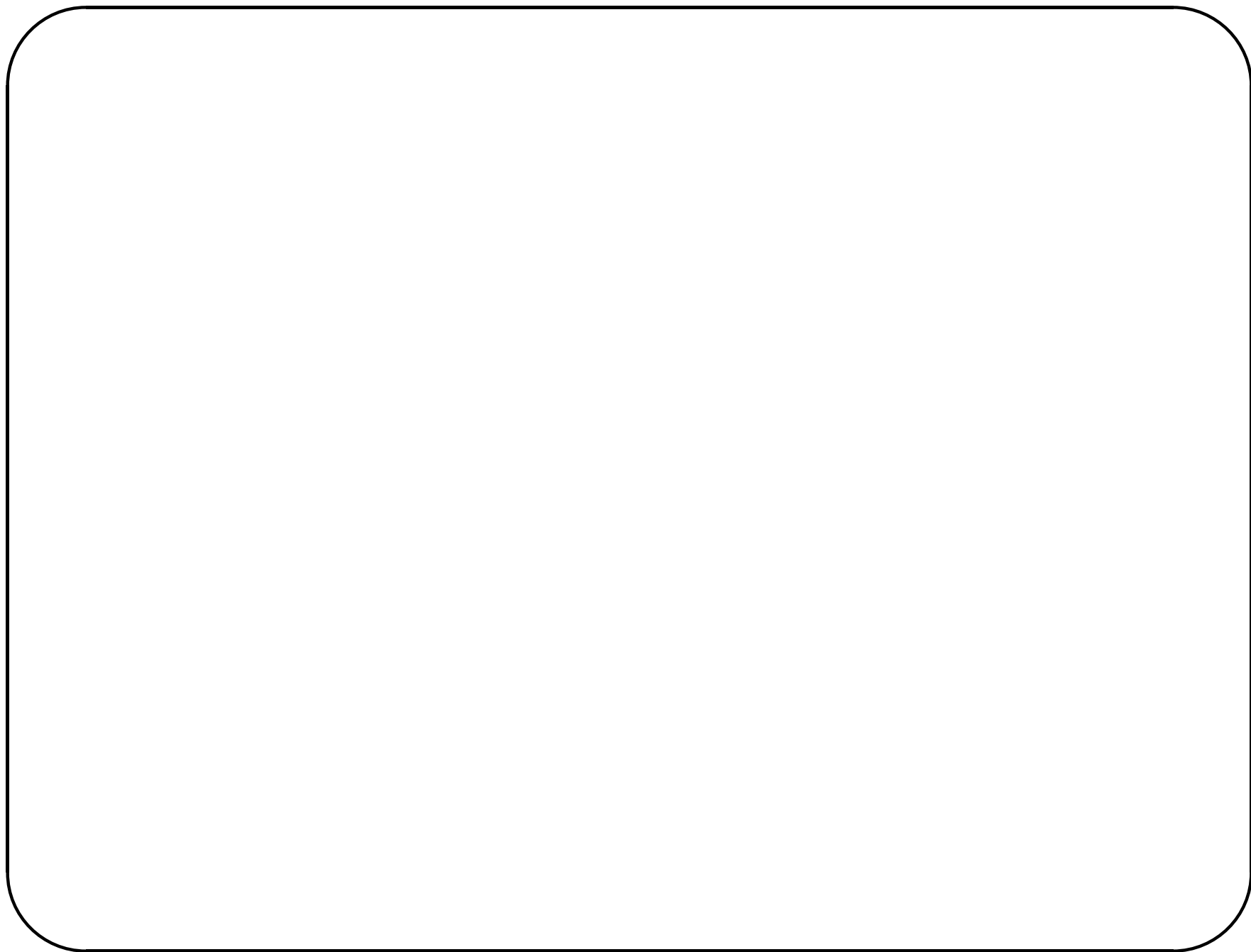
^aPlasma Confinement, R.D. Hazeltine, and J.D. Meiss (Dover, 2003).

Slab Approximation



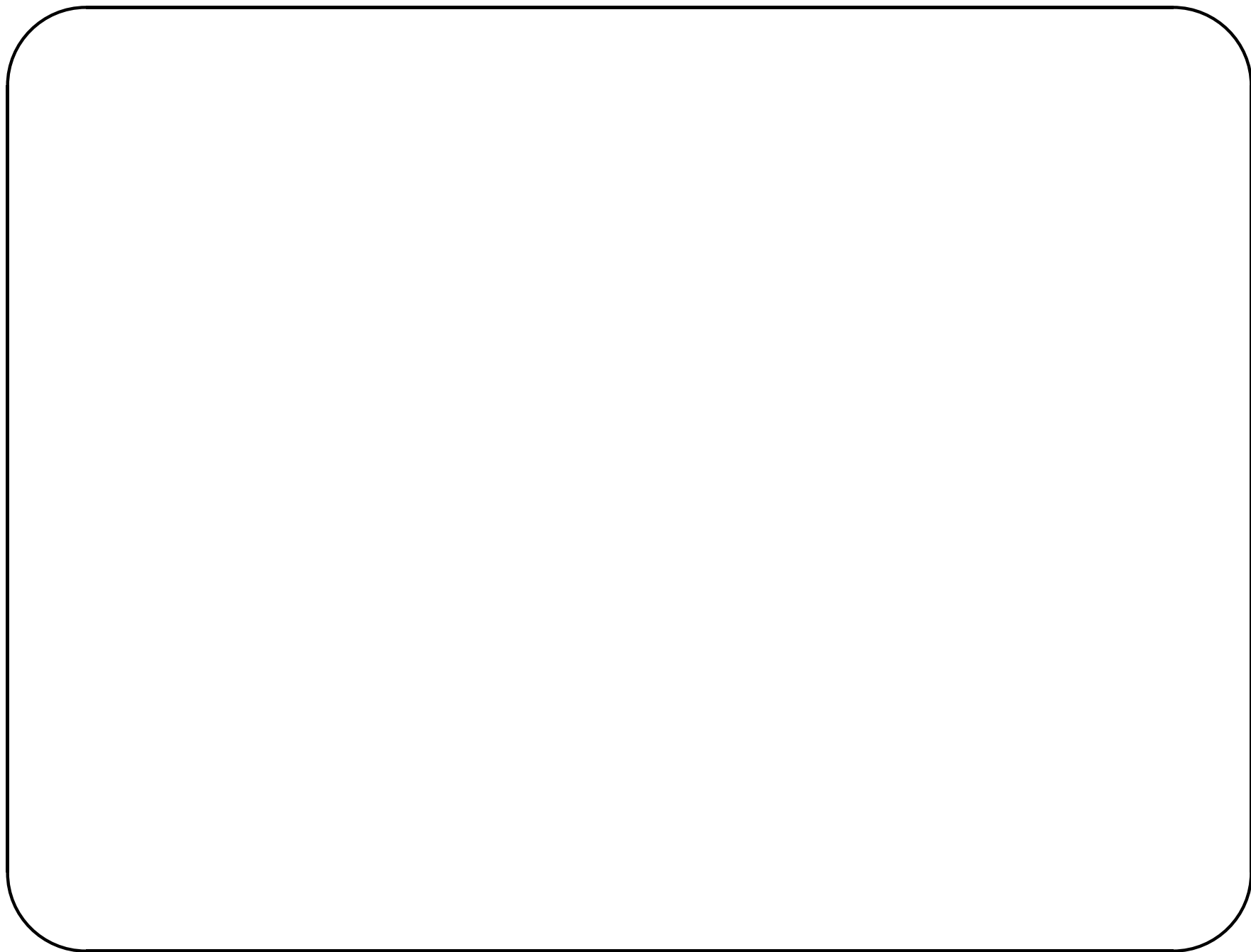
Slab Model

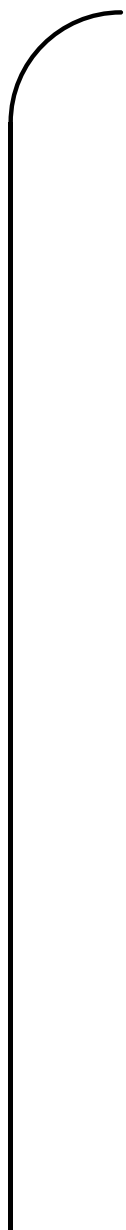
- Cartesian coordinates x, y, z $L \leq z \leq 0$
 - Assumption: field in z direction $\vec{B} = B_z \vec{z}$
 - Ampere's law $\nabla \times \vec{H} = \vec{j}$ $\Rightarrow \frac{dB_z}{dz} = -\mu_0 j_y$
 - Ampere's law $\nabla \times \vec{H} = \vec{j}$ $\Rightarrow \frac{dB_z}{dz} = -\mu_0 j_y$
- $$L_s = B_z / B'_z(0).$$
- Ampere's law $\nabla \times \vec{H} = \vec{j}$ $\Rightarrow \frac{dB_z}{dz} = -\mu_0 j_y$
 - Perfect conductor boundary conditions $x = \pm a$
 - Normal component of \vec{H} is zero $\vec{k} = (0, k, 0)$ $\Rightarrow \vec{k} \cdot \vec{B} = 0$ at $x = 0$



Outer Region

- n t r r g n f f p r s s s t f p s n n n r



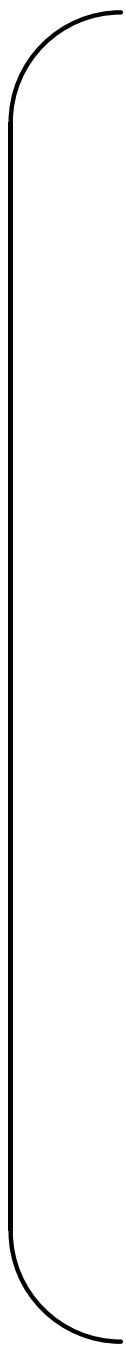


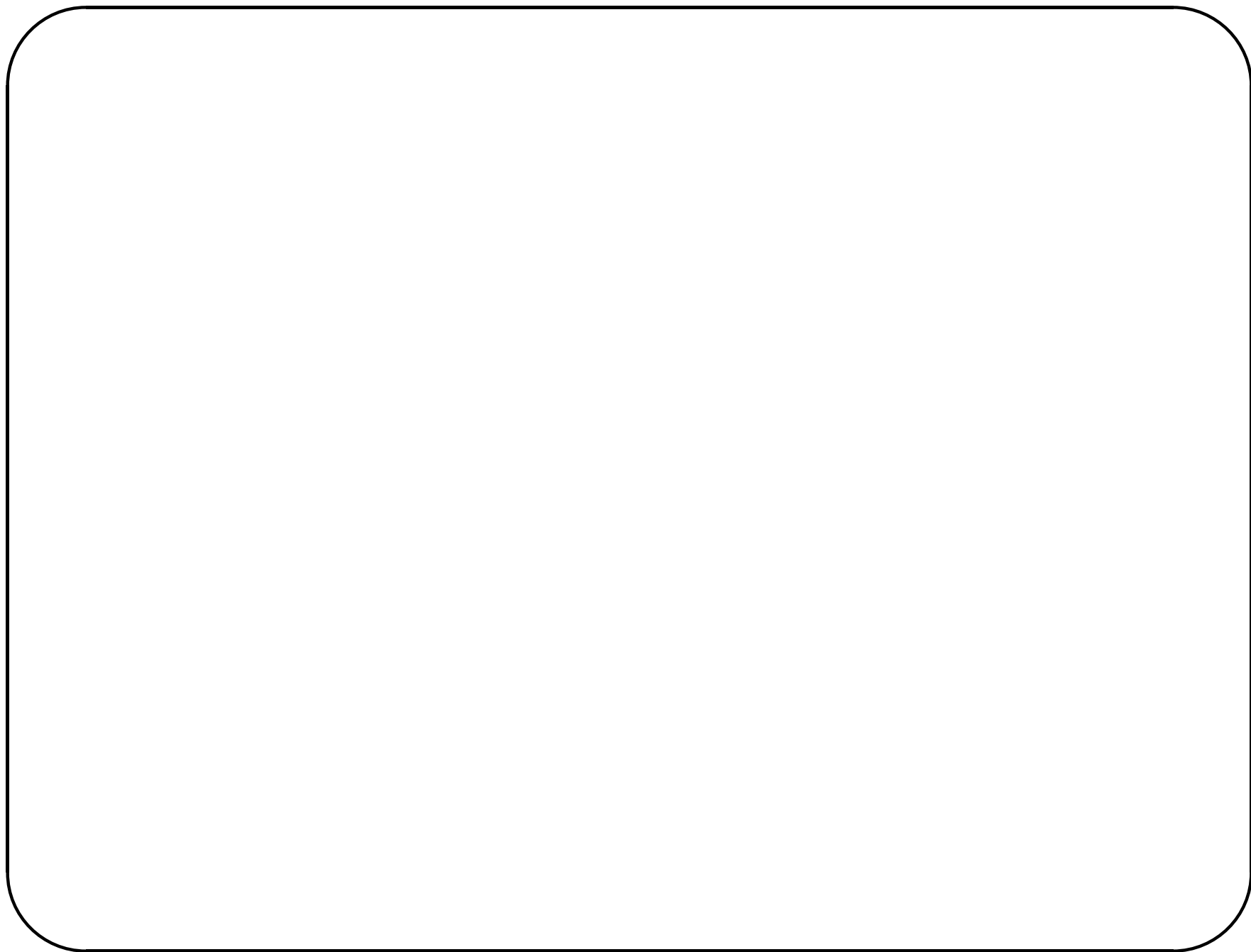
Inner Region

- $\frac{1}{2} \int_{-1}^1 \frac{1}{\sqrt{1-x^2}} dx = \frac{1}{2} \left[\arcsin x \right]_{-1}^1 = \frac{1}{2} (\frac{\pi}{2} - (-\frac{\pi}{2})) = \frac{\pi}{2}$
- $\int_0^1 x^2 dx = \left[\frac{x^3}{3} \right]_0^1 = \frac{1}{3}$
- $\int_0^1 x^2 dx = \frac{1}{3}$

Constant- Approximation

- $(1)(X)$ g n r y σ s n t r y s g n c n y n X r n n r r g n
 $| (1)(W) - (1)(0) | | (1)(0) |$
- Constant- approximation* t r t $(1)(X)$ s c n s t n n X





MHD Flow - II

- L_t

$$M(x) = \frac{d}{dx}.$$

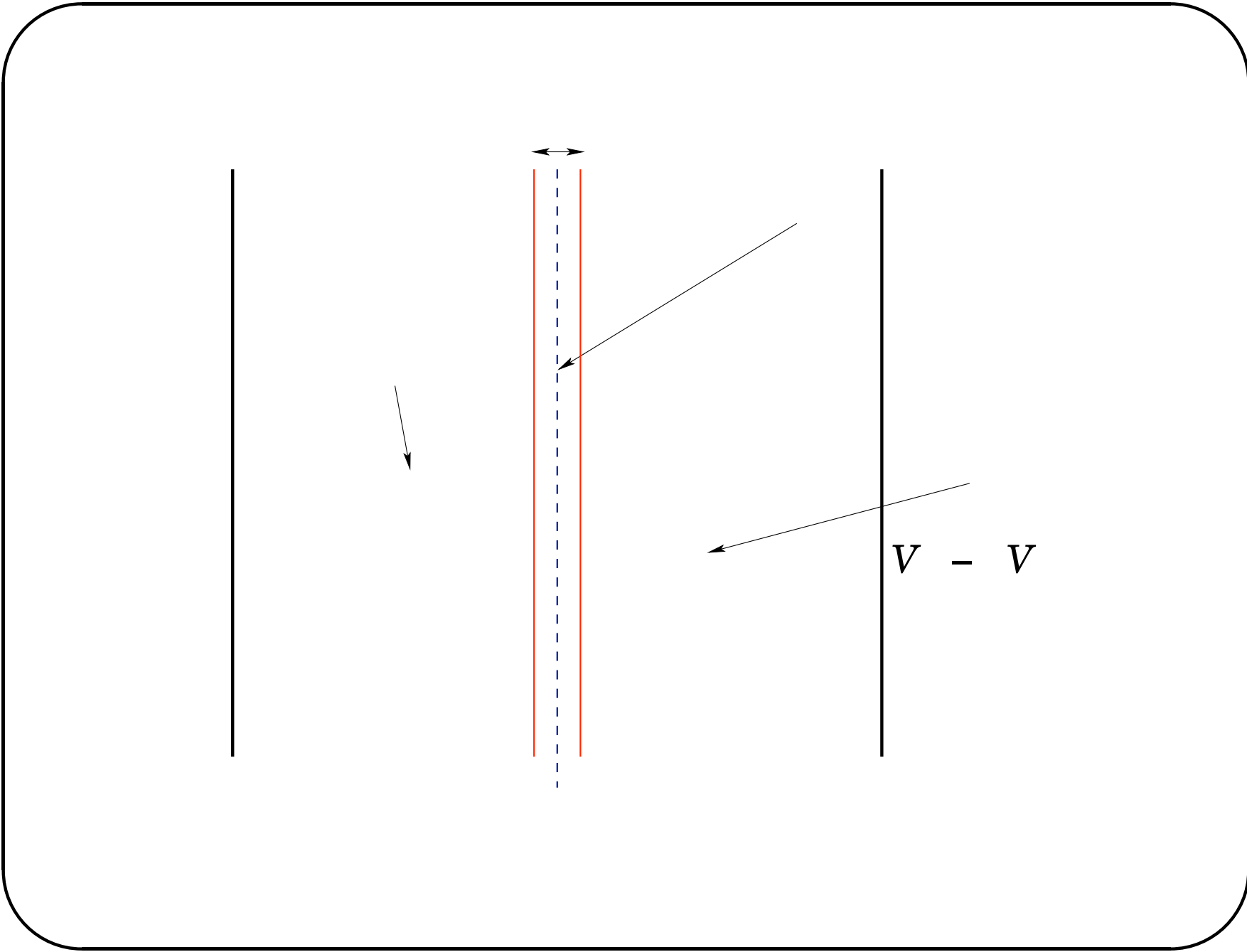
- $\psi = y^2 M(x)$

$$V_y = x M(x).$$

- $\psi = y^2 M(x)$ is *odd* function of x

$$M = 0$$

$\psi = y^2 M(x)$ is *odd* function of x is *odd* function of x
 $\psi = y^2 M(x)$ is *odd* function of x is *odd* function of x



Rutherford Equation - I

- Asymptotic expansion of the scattering amplitude $f(\theta)$ for large k

$$f(\theta) \sim -\frac{1}{k} \int_{-\infty}^{+\infty} J(\alpha) \cos(\alpha \theta) d\alpha.$$

- The function $J(\alpha)$ is the Fourier transform of the scattering potential $V(r)$

$$J(\alpha) = \int_0^{\infty} V(r) r \cos(\alpha r) dr.$$

- The function $J(\alpha)$ is

$$J(\alpha) = J(\alpha).$$

Rutherford Equation - II

- O_h s

$$\frac{d}{dt} \cos$$

Rutherford Equation - IV

- $\frac{dW}{dt} = 0.41 \left(- \frac{d^4 B_y^{(0)} / dx^4}{d^2 B_y^{(0)} / dx^2} \right)_{x=0} W.$

$$\frac{dW}{dt} = 0.41 \left(- \frac{d^4 B_y^{(0)} / dx^4}{d^2 B_y^{(0)} / dx^2} \right)_{x=0} W.$$

- $\frac{dW}{dt} = 0$

W

MHD Theory: Summary

- ring α inst $f' > 0$
-

Drift-MHD Theory

- n_e drift $\nabla \cdot \mathbf{p}_p$ \rightarrow $\mathbf{v}_E = \frac{c}{4\pi n_e} \nabla \times \mathbf{B}$ \rightarrow *charged particle drift velocities*
- $\nabla \cdot \mathbf{v}_E = 0$ \rightarrow *theory*
- $\nabla \cdot \mathbf{v}_E = 0$ \rightarrow *ion Larmor radius calculated*
- $\nabla \cdot \mathbf{v}_E = 0$ \rightarrow $\mathbf{v}_* = \frac{c}{4\pi n_e} \nabla \times \mathbf{B}$

Basic Assumptions

- Retains σ frequency specificity
- Assumes perfect horizontal transport capacity
 $T_e = T_e(\quad)$
- Assumes $T_i/T_e = \quad = \text{constant}$ frequency specificity

Basic Definitions

- r_s
 - $\frac{g_{nc}}{f_{atn}}$
 - J_{prcrrn}
 - $\frac{g_{ngc}}{r_{e}} \frac{str}{f_{atn}}$
 - $U_{prnrc ty}$
 - $n_{ctrnn} r_{nsty} n_{snfr} c_{grn}$
 - V_z prn cty
- $P_{rt rs}$
 - $= (L_n/L_s)^2 \frac{h r L_n s q}{r_{nsty} gr_{n}}$
 - $\frac{r_{sst ty} D}{p r p_{nc} r p r c_{s ty} \mu_i/e}$
 - $\frac{p r p_{nc} r n/ctrn s sty}{}$

Drift-MHD Equations - I

- ψ system ψ drift $\nabla \psi \cdot \nabla \psi = -q \psi^2$

$$\psi = -x^2/2 + \cos \theta, \quad U = \psi^2,$$

$$0 = [\psi, n] + J_z,$$

$$0 = [\psi, U] - \frac{1}{2} \{ \psi^2 [\psi, n] + [U, n] + [\psi^2 n, \psi] \}$$

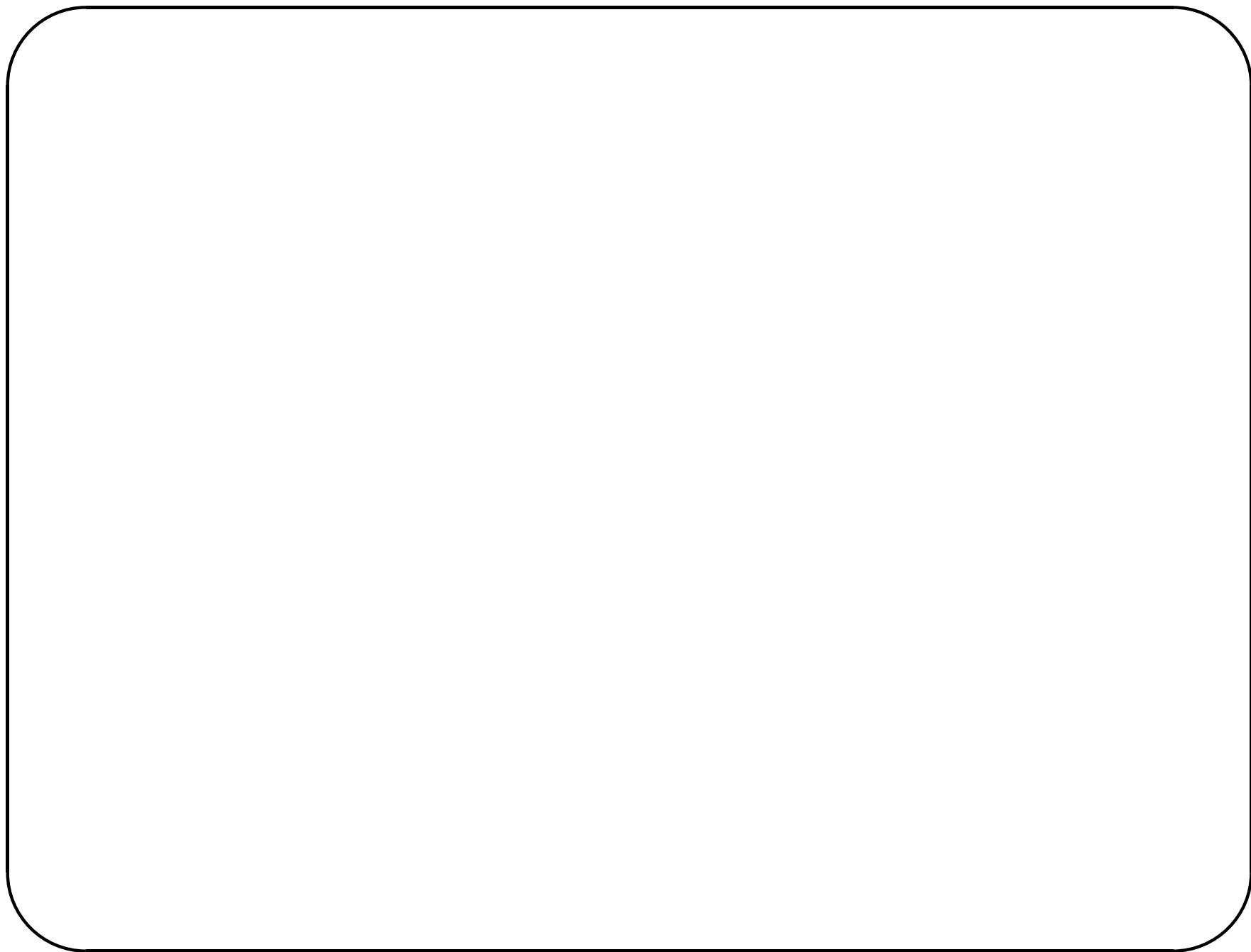
$$+ [J_z, \psi] + \mu_i \nabla^2 (\psi + n) + \mu_e \nabla^2 (\psi - n),$$

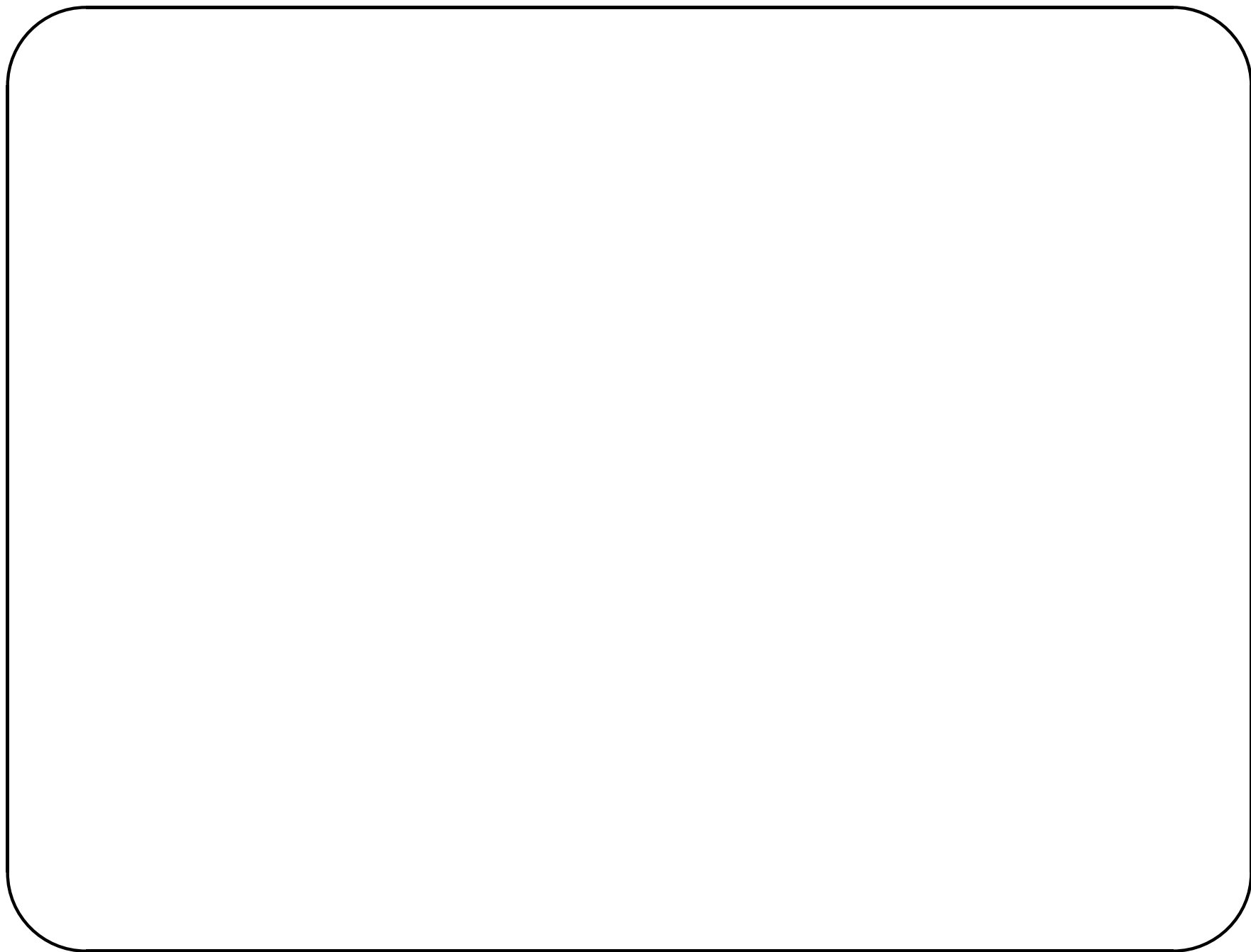
$$0 = [\psi, n] + [V_z + J_z, \psi] + D \nabla^2 n,$$

$$0 = [\psi, V_z] + [\psi, J_z]$$

Drift-MHD Equations - II

- y try





Subsonic Islands^a

- $$s_{t r} = \dots = (\dots), n = n(\dots).$$
- $$s_{t r} = \dots = -n \dots$$

Analysis - I


 ns ty q t nr c st

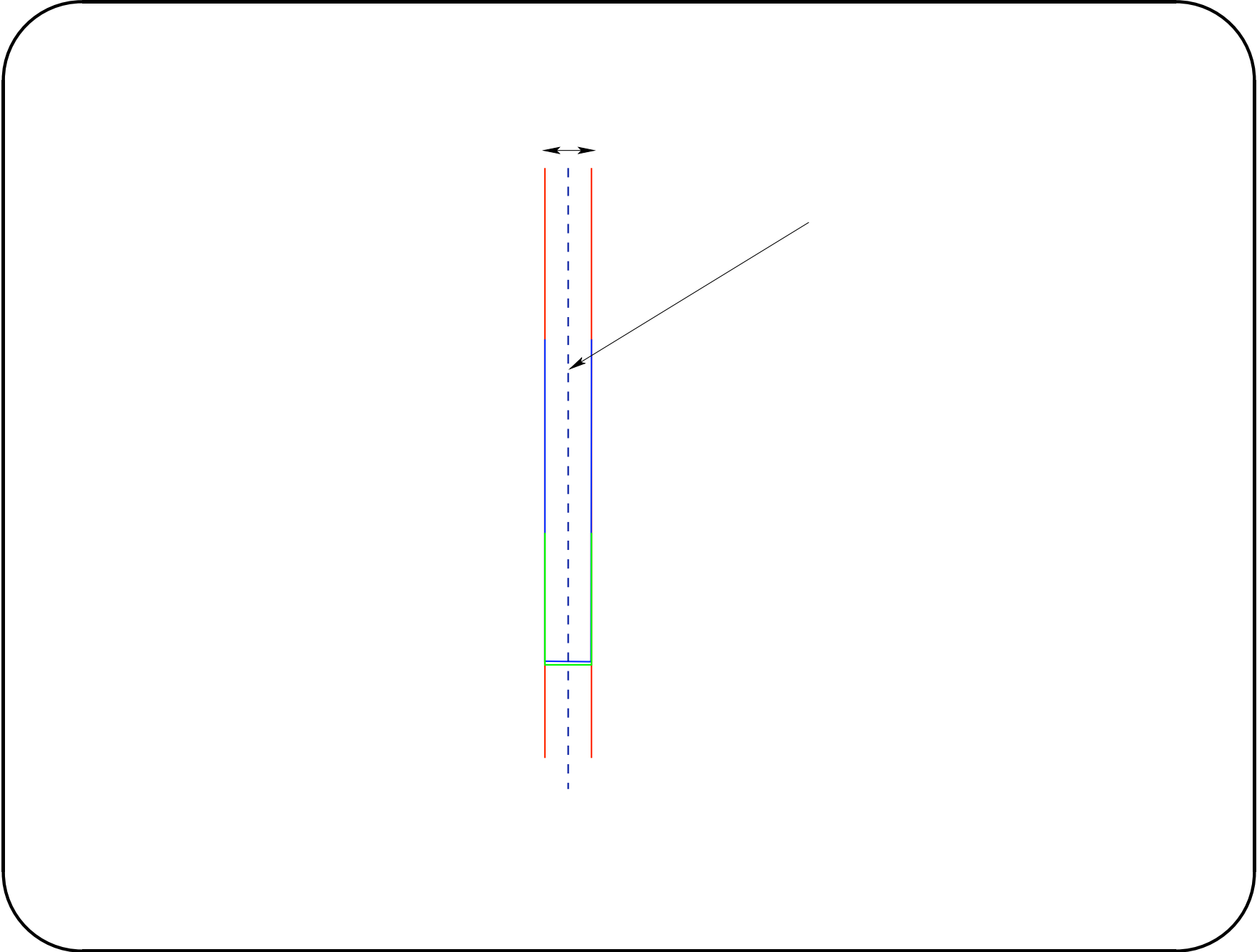
$$0 [V_z + J,] + D^2 n.$$

• rc ty q t nr c st

$$0 [-MU - (\ /2)(LU + M^2 n) + J,]$$

$$+ \mu_i^4 (+ n) + \mu_e^4 (- n).$$

• s rfc rg th q t ns rc ng th t [A,] = 0



Island Propagation

- As $|x|/W \rightarrow \infty$, $v_{yE \times B} \rightarrow V_{EB} - V_{h r}$ V_{EB} s

-

$$V = V_{EB} + \frac{(\mu_i - \mu_e)}{(1 +)(\mu_i + \mu_e)}.$$

-

$$V_i = V_{EB} + 1/(1 +), \quad V_e = V_{EB} - 1/(1 +).$$

-

$$V = \frac{\mu_i}{\mu_i + \mu_e} V_i + \frac{\mu_e}{\mu_i + \mu_e} V_e.$$

scos/y elon ed a e a e f

Polarization Term - I

- $\tau_c \propto \tau_q \propto \tau_{hy} \propto \tau^s$

$$J_c = \frac{1}{2} \left(x^2 - \frac{x^2}{1} \right) \frac{d[M(M+L)]}{d} + I(\)$$

$\tau^s \propto \tau_{spr} \propto \tau_{hr} \propto J_c \propto \tau_{spr} \propto \tau_{fh} \propto \tau_{h} \cos \theta \propto \tau_{try}$

- As $\tau_{spr} \propto \tau_{rfc} \propto \tau_{rg} \propto \tau_{oh} \propto \tau_{sy} \propto \tau^s$

$$J_c = I(\) \propto \tau^s = -1 \frac{d}{dt} \cos \theta$$

- τ^s

$$J_c = \frac{1}{2} \left(x^2 - \frac{x^2}{1} \right) \frac{d[M(M+L)]}{d} + -1 \frac{d}{dt} \frac{\cos \theta}{1}$$

Polarization Term - II

- Asymptotic expansion of the polarization term

$$P_1 = -4 \int_0^\infty J_0(\alpha r) \cos(\alpha z) d\alpha$$

- Asymptotic expansion of the polarization term

Drift-MHD Theory: Summary

- Resistivity is neglected in the Ohmic law, but it is important in the induction equation.
- The parallel component of the induction equation is decoupled from the perpendicular components.
- The parallel component of the induction equation is a diffusion equation for the parallel magnetic field.
- The perpendicular components of the induction equation are coupled to the parallel component through the drift velocity.