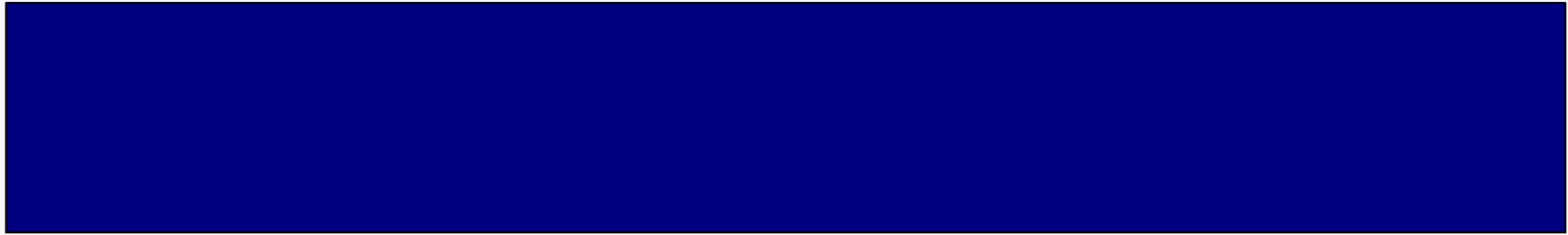


Plasma Rotation in



Mechanisms have been developed to control/affect plasma rotation

Describing toroidal rotation in tokamaks
is a transport problem

Thesis

Outline

Starting point for the calculation is
the plasma kinetic equation

(x, v)

$$\frac{df_s}{dt} + \nabla_{\vec{v}} \cdot$$

$$d^3\vec{v} \left(1, m_s \vec{v}, \frac{m_s v^2}{2} \right) \left[\frac{\partial f_s}{\partial t} + \vec{v} \cdot \nabla f_s + \frac{q_s}{m_s} (\vec{E} + \vec{v} \times \vec{B}) \cdot \nabla_{\vec{v}} f_s = C(f_s) + S(f_s) \right]$$

V_s

A number of assumptions are made
to make analytic progress

Small gyroradius expansion is used

Transport equations for density and pressure are obtained by flux surface averaging

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The different orders of the momentum balance equation refer to different timescales

$$\vec{J}_0 \quad \vec{B}_0 = \# p_0$$

$$= e \quad n_i q_i E + V_i \quad B \quad p_i$$

$$\Omega_i \quad V_i \quad V = \frac{d}{d} + \frac{dp_i}{n_i q_i d} \quad qV$$

$$V_i \quad \frac{E_r}{B_p} \quad \frac{dp_i}{n_i q_i dr} + \frac{B_t}{B_p} V_p$$

∇

$$\frac{r}{V} = e \frac{r}{V\#} \quad + e \frac{r}{\%} \frac{r}{V\#} \quad \% = V \frac{\dot{B}}{B} + \frac{r}{V\&}$$

Poloidal flow is obtained from parallel force balance

Viscous damping occurs

After determining the poloidal flow, there is a unique relationship between E_r and the toroidal rotation

$$\Omega \equiv \vec{V} \cdot \nabla \zeta = -\left(\frac{d\Phi}{d\psi} + \frac{1}{n_i q_i} \frac{dp_i}{d\psi} - q \vec{V} \cdot \nabla \theta \right)$$

$$\langle \vec{B} \cdot \vec{v}_i \rangle \%$$

$$V_t \frac{E_r}{B_p} \# \frac{1}{n_i q_i B_p} \frac{dp_i}{dr} + \frac{1.17}{q_i B_p} \frac{dT_i}{dr}$$

Pressure/temperature profiles
determined by transport processes

$$\begin{aligned} \frac{\vec{r}}{V} = \frac{\vec{r}}{e} \frac{\vec{r}}{V} + \frac{\vec{r}}{e} \frac{\vec{r}}{V} &= V \frac{\vec{B}}{B} + \frac{\vec{r}}{V} \quad \frac{\vec{r}}{V} = \frac{\vec{B}}{B} \frac{d}{d} + \frac{dp_s}{n_s q_s d} \\ V & R \frac{d}{d} + \frac{dp_i}{n_i q_i d} + \frac{R}{q_i R} \frac{dT_i}{d} \end{aligned}$$

Electron parallel momentum balance produces parallel Ohm's law

$$e \vec{E} \cdot \vec{B} = - \vec{B} \cdot \nabla_{\parallel} \tilde{\phi}_e - \vec{B} \cdot \vec{R}_e - \vec{B} \cdot \vec{S}_e$$

$$e \vec{B} \cdot \vec{V}_e = - \nabla_{\parallel} \tilde{\phi}_e - \vec{B} \cdot \vec{V}_e \cdot \vec{B}$$

#

$$\langle \tilde{\phi}_e \rangle = \langle \tilde{\phi}_e \rangle + \mu$$

$$\vec{E} \cdot \vec{B}_0 = - \nabla_{\parallel} \tilde{\phi}_e - \vec{B}_0 \cdot \nabla_{\parallel} [\vec{J} \cdot \vec{B}_{BS} + \vec{J} \cdot \vec{B}_{CD} + \vec{J} \cdot \vec{B}_{dn}]$$

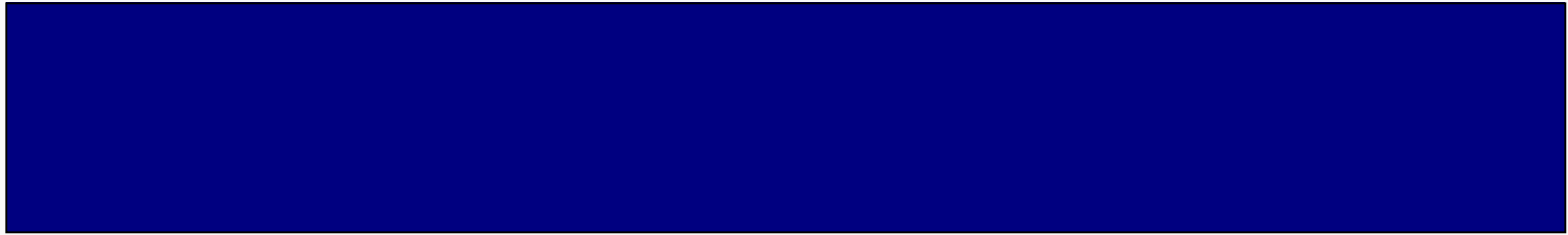
$$= - \frac{\mu}{\tau} \langle \tilde{\phi}_e \rangle + \mu \langle \tilde{\phi}_e \rangle - \mu \langle \tilde{\phi}_e \rangle$$

$$\langle \tilde{\phi}_e \rangle = \mu \langle \tilde{\phi}_e \rangle$$

Parallel Ohm's law is used to describe collisional ambipolar particle flux

$$\frac{a}{s} = \frac{1}{q_s} \langle e \frac{r}{R_s} \rangle - n_0 \langle e \frac{r}{E} \rangle \quad \frac{r}{R_e} n_e e (\frac{r}{J_{\parallel}} + \frac{r}{J}) = \frac{r}{R_i}$$

Plasma fluctuations influence particle flux/toroidal
momentum balance

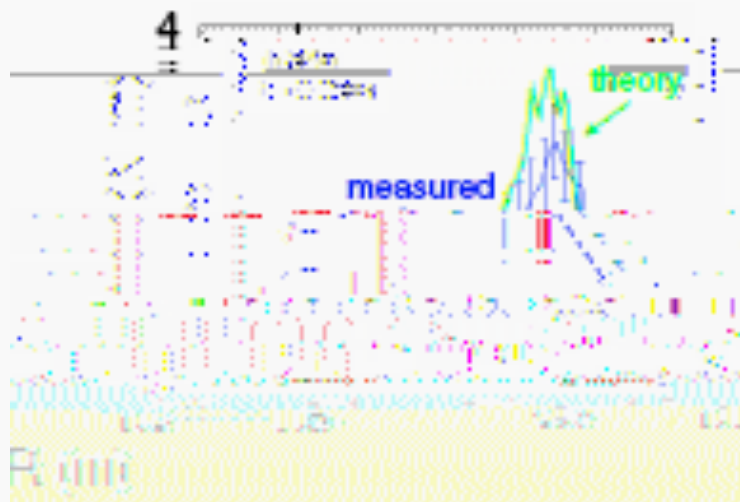


3-D magnetic fields produce neoclassical toroidal viscous forces (NTV) throughout the plasma

μ

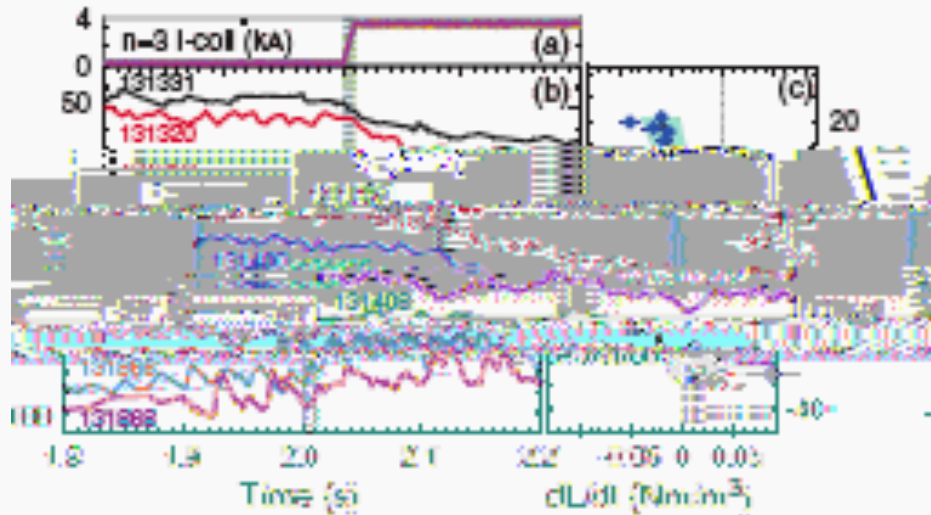
μ

The NTV force is felt throughout the plasma



- Favorable comparison to analytic predictions
- NTV physics has been seen on NSTX, DIII-D, MAST, JET

Experiments on DIII-D have demonstrated the presence of the NTV offset velocity



Initially, slowly rotating
Plasmas sped up to the
Offset NTV velocity when
3-D fields are applied

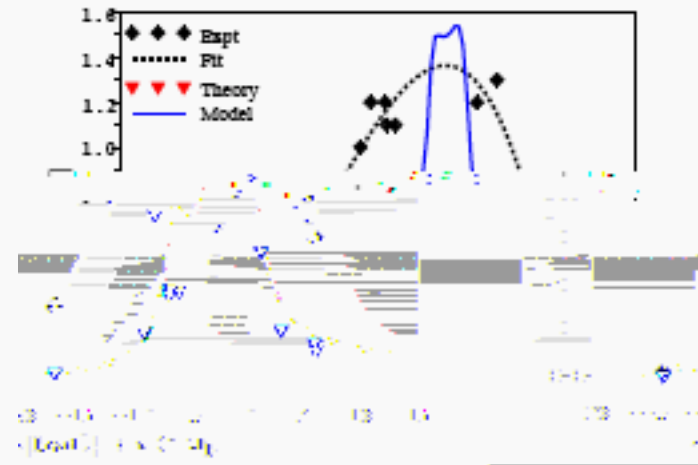
Recent experiments on DIII-D have demonstrated a peak in the NTV force at zero radial electric field

μ

$$i \frac{R}{\%B} \frac{nv_{ii} \sqrt{\#}}{\sqrt{\#}} \nu E$$

Peaks at

Recent experiment on DIII-D
Demonstrates peak NTV at



Particle flux has 8 non-ambipolar contributions

Zero radial current produces torque balance relation

$$\langle \mathbf{j} \cdot \nabla \psi \rangle = \sum_s q_s \Gamma_s = \sum_s q_s \Gamma_s^{NA}$$

$$\frac{\partial}{\partial t} \langle \rho \rangle_q$$

Toroidal rotation equation includes many different effects

Ω

$$\frac{1}{V'} \frac{d}{dt} (V' L) + \vec{e}_r \cdot \nabla_{\parallel}^{NA} + \vec{e}_r \cdot \nabla_{\perp} \left(\frac{1}{V'} \frac{d}{dt} (V' \cdot) \right) + \vec{e}_r \cdot (\tilde{\mathbf{J}} + \tilde{\mathbf{B}}) + \frac{L}{V'} \vec{e}_r \cdot \vec{S}$$

$<^r$

Toroidal rotation determines radial electric field
required for net ambipolar particle flux

Summary



Summary